Proceedings of the 15th Czech and Slovak Conference on Magnetism, Košice, Slovakia, June 17-21 2013

# Magnetism in $RCo_2$ (R = Dy, Ho, Er, Tm) Under Hydrostatic Pressure

J. Prchal\*, J. Šebesta, J. Valenta, M. Míšek, D. Turčinková, L. Lapčák,

J. Prokleška, M. Kratochvílová and V. Sechovský

Charles University in Prague, Faculty of Mathematics and Physics,

DCMP, Ke Karlovu 5, 121 16 Prague 2, Czech Republic

The short-range parimagnetic configurations, observed in the heavy rare-earth RCo<sub>2</sub> compounds in paramagnetic range far above the Curie temperature  $T_C$ , consist in formation of ferromagnetic cobalt clusters antiferromagnetically coupled to the neighboring R magnetic moments. The characteristic temperature of the onset of parimagnetism  $T_f$  is very sensitive to changes of external conditions like pressure or composition. The pressure coefficients  $\partial T_f/\partial p$  and  $\partial T_C/\partial p$  are comparable, indicating a close connection of underlying mechanisms. Our recent measurements indicate systematic evolution of the pressure coefficients with spanning the rare-earth series from Dy through Tm in RCo<sub>2</sub> with a specific case TmCo<sub>2</sub>.

DOI: 10.12693/APhysPolA.126.288

PACS: 75.20.En; 75.50.Gg

#### 1. Introduction

The so called parimagnetism, which was discovered few years ago [1], appearing in heavy rare-earth RCo<sub>2</sub> compounds (R = Dy, Ho, Er, Tm) makes the study of compounds from this family attractive again [1, 2]. In the  $RCo_2$  (R = Dy, Ho, Er, Tm) compounds crystallizing in the cubic Laves phase structure the localized rare-earth 4f-electron magnetic moments coexist together with the itinerant Co 3d moments, which appear on the verge of magnetism. Below  $T_C$ , the large exchange field due to the ferromagnetically ordered R moments polarizes the Co 3d-electron states and the emerged moments in the Co sublattice orient antiparallel to the R sublattice. Recent reports propose that weak Co moments survive rather far above  $T_C$  in Co ferromagnetic clusters, remaining coupled antiparallel to paramagnetic R moments forming a parimagnetic configuration in the paramagnetic phase. The onset of this phenomenon causes an anomaly in the AC-susceptibility at  $T_f$  – the "characteristic temperature" of the onset of parimagnetism".

We have performed a systematic study of pressure influence on both the Curie temperature  $T_C$  and the parimagnetic temperature  $T_f$  in several RECo<sub>2</sub> compounds with R = Dy, Ho, Er and Tm. Whereas the detailed results of these high-pressure experiments will be presented elsewhere, here we present summary of the performed measurements.

# 2. Experimental

Polycrystalline samples were prepared by arc-melting of the stoichiometric ratio of the constituent elements. Each sample was turned and remelted several times for better homogeneity. After annealing and tempering the samples were checked for their quality by X-ray diffraction and the real composition was determined by SEM

with the EDX analyzer. Small amount (less than 3%) of RCo<sub>3</sub> impurity phase was found in the samples, except for TmCo<sub>2</sub>. For measurements of bulk properties (AC-susceptibility, electric resistivity, specific heat) a commercial instruments PPMS, MPMS (both Quantum Design) and Closed Cycle Cryocooler (CCR; Janis Research) were used. Experiments under hydrostatic pressure up to 3 GPa were performed using double-layered clamped pressure cell made of Cu-Be bronze and Ni-Cr-Al alloy with special design to fit into both – PPMS and CCR. Daphne 7373 oil [3] was used as a pressuretransmitting medium and a pressure dependence of resistivity of a manganin wire was used to determine the pressure inside the pressure cell. For the experiment of magnetic susceptibility under pressure a homemade miniature detection coil set was used, which fits inside the bore of the pressure cell [4].

## 3. Results and discussion

In AC-susceptibility data for all the studied samples, both the Curie temperature anomaly and the much weaker anomaly at the parimagnetic temperature  $T_f$  are clearly visible (see Table). The determination of  $T_C$  was supported by the corresponding phase-transition effects present in specific heat and electrical resistivity data, collected at ambient pressure. The estimated Curie temperature values are in agreement with those published before [4–8].

TABLE Curie-  $(T_C)$  and parimagnetic  $(T_f)$  temperatures of the studied RCo<sub>2</sub> compounds at ambient pressure.

	$T_C$ (K)	$T_f$ (K)
$\mathrm{DyCo}_2$	137	200
$\mathrm{HoCo}_2$	80	125
${ m ErCo_2}$	33	98
${\rm TmCo_2}$	3.7	35

 $<sup>\</sup>hbox{$^*$ corresponding author; e-mail: $\tt prchal@karlov.mff.cuni.cz$}$ 

Apparently the ordering temperature appears to be very sensitive to applied hydrostatic pressure. Comparable values of pressure coefficient has been found also for  $T_f$  for corresponding compounds, which indicates a common underlying mechanism of ferrimagnetism and parimagnetism (see Fig. 1). Both,  $T_C$  and  $T_f$ , shift to lower temperatures with pressure, exhibiting a monotonous evolution of the pressure coefficients across the rare-earth series. The last compound TmCo<sub>2</sub> seems then to be a rather exceptional case with the pressure coefficient being almost zero. The pressure coefficients approaching the zero with the increase of atomic number of the rare earth can be connected with the decrease of compressibility of the respective compound, which unfortunately could not be measured experimentally, but it can be expected with respect to lanthanide contraction causing a reduction of the interatomic space. Thus also the strength of exchange interaction will be less influenced by the external pressure, which is exactly what we observe.

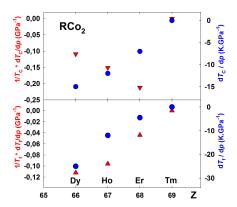


Fig. 1. Evolution of the pressure coefficients and normalized pressure coefficients of the characteristic temperatures  $T_C$  and  $T_f$  for RCo<sub>2</sub> with R = Dy, Ho, Er, Tm. The values for HoCo<sub>2</sub> and ErCo<sub>2</sub> were taken from refs. [8] and [4], respectively.

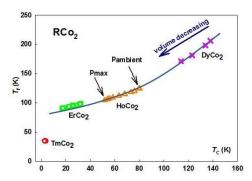


Fig. 2. Empirically estimated relation between the characteristic temperatures  $T_C$  and  $T_f$  in the studied RCo<sub>2</sub> compounds under various pressures.  $P_{max} = 1.5$  GPa for R = Dy, whereas  $P_{max} = 3$  GPa for R = Tm, Er and Ho.

When searching for a more general relation between  $T_C$  and  $T_f$ , we have estimated a monotonous and rather smooth evolution, as can be seen in Fig. 2. Here we can see a clearly deviating position of all the pressure points of  ${\rm TmCo_2}$  from the intimated dependence. Although we cannot fully explain the origin of the mentioned behavior yet, it is already clear that the  $T_C$  vs.  $T_f$  dependence is systematic with the decrease of the unit-cell volume, and that  ${\rm TmCo_2}$  is excess compound with practically zero dependence of  $T_C$  and  $T_f$  on the applied pressure.

### 4. Conclusions

We may conclude that our findings of evolution of the Curie temperatures and of the parimagnetic temperatures with the application of external pressure are always characterized by a negative pressure coefficient, which in most of the studied compounds is relatively high in absolute values and systematically evolves when spanning the rare-earth series from Dy through Tm. The exception is  ${\rm TmCo_2}$  which exhibits almost no change of  $T_C$  and  $T_f$  with pressure up to 3 GPa. This also makes  ${\rm TmCo_2}$  an unusual case in the proposed general evolution  $T_C$  vs.  $T_f$  for other  ${\rm RCo_2}$  compounds. Further study for explanation of the pressure influence on the magnetism in  ${\rm RCo_2}$  (TmCo<sub>2</sub>) is necessary.

## Acknowledgments

The experiments were performed in MLTL (http://mltl.eu/), which is supported within the program of Czech Research Infrastructures (project no. LM2011025). The work was also supported by the Czech Science Foundation (project No. P204/12/0692).

#### References

- [1] J. Herrero-Albillos, F. Bartolomé, L.M. García, A.T. Young, T. Funk, J. Campo, G.J. Cuello, *Phys. Rev. B* 76, 094409 (2007).
- [2] C.M. Bonilla, I. Calvo, J. Herrero-Albillos, A.I. Figueroa, C. Castan-Guerrero1, J. BartolomĂŠ1, J.A. Rodriguez-Velamazan, D. Schmitz, E. Weschke, D. Paudyal, V.K. Pecharsky, K.A. Gschneidner Jr., F. BartolomĂŠ1, L.M. GarcĂa1, J. Appl. Phys. 111, 07E315 (2012).
- [3] K. Murata, H. Yoshino, H.O. Yadav, Y. Honda, N. Shirakawa, Rev. Sci. Instrum. 68, 2490 (1997).
- [4] M. Míšek, J. Prokleška, V. Sechovský, D. Turčinková, J. Prchal, A.F. Kusmartseva, K.V. Kamenev, J. Kamarád, J. Appl. Phys. 111, 07E132 (2012).
- [5] E. Gratz, R. Resel, A.T. Burkov, E. Bauer, A.S. Markosyan, A. Galatanu, J. Phys.: Cond. Matter. 7, 6687 (1995).
- [6] N.K. Singh, K.G. Suresh, A.K. Nigam, Sol. Stat. Commun. 127, 373 (2003).
- [7] R. Resel, M. Hedo, T. Nakama, K. Yagasaki, E. Gratz, R. Hauser, A.S. Markosyan, Sol. Stat. Commun. 95, 735 (1995).
- [8] J. Valenta, J. Prchal, M. Kratochvílová, Acta Phys. Pol. A 126, 404 (2014).