

Supplementary material

Non-linearity coefficient μ

1. The variation with pressure of the volume of a system for which μ is constant is given by $V(P) = V_0(1 + aP)^{-\frac{1}{\mu}}$ where V_0 is a constant, which represents the volume of the system in the limit of vanishing pressure.

Since $\mu = -\left[\frac{\partial}{\partial P}\left(V/\left(\frac{\partial V}{\partial P}\right)_T\right)\right]_T = V\left(\frac{\partial^2 V}{\partial P^2}\right)_T / \left(\frac{\partial V}{\partial P}\right)_T^2 - 1$, an expression for $\mu(P)$ can easily be

derived when the variation with pressure of the system volume is known as a power series

$$\begin{aligned} V(P) &= \sum_{i=0}^{\infty} a_i P^i : \quad \mu(P) = \sum_{i=0}^{\infty} a_i P^i \sum_{j=0}^{\infty} (j+1)(j+2)a_{j+2} P^j \bigg/ \sum_{i=0}^{\infty} (i+1)a_{i+1} P^i \sum_{j=0}^{\infty} (j+1)a_{j+1} P^j - 1 \\ &= \sum_{k=0}^{\infty} P^k \sum_{l=0}^k (l+1)(l+2)a_{k-l} a_{l+2} \bigg/ \sum_{k=0}^{\infty} P^k \sum_{l=0}^k (l+1)(k-l+1)a_{l+1} a_{k-l+1} - 1. \end{aligned}$$

In particular, $\mu = 2\frac{a_0 a_2}{a_1^2} - 1$ in the limit of vanishing pressure.

2. The Lennard-Jones potential can be written as $E_p(r) = E_0 \left[\left(\frac{r_e}{r}\right)^{12} - 2\left(\frac{r_e}{r}\right)^6 \right]$, where E_0 is a positive constant and r_e is an equilibrium distance for which

$$E_p(r_e) = E_0 \left[\left(\frac{r_e}{r}\right)^{12} - 2\left(\frac{r_e}{r}\right)^6 \right]_{r=r_e} = -E_0,$$

$$\frac{dE_p}{dr}(r_e) = -\frac{12E_0}{r_e} \left[\left(\frac{r_e}{r}\right)^{13} - \left(\frac{r_e}{r}\right)^7 \right]_{r=r_e} = 0,$$

$$\frac{d^2 E_p}{dr^2}(r_e) = \frac{12E_0}{r_e^2} \left[13\left(\frac{r_e}{r}\right)^{14} - 7\left(\frac{r_e}{r}\right)^8 \right]_{r=r_e} = \frac{72E_0}{r_e^2} > 0, \text{ which means that for } r = r_e, E_p(r) \text{ is}$$

minimal.

Assume that the volume of the system is proportional to the cube of the distance r , which

translates into $V(r) = V_e \left(\frac{r}{r_e}\right)^3$, then:

$$E_p(r) = E_0 \left[\left(\frac{V_e}{V} \right)^4 - 2 \left(\frac{V_e}{V} \right)^2 \right] \text{ and}$$

$$P(V) = - \frac{\partial E_p}{\partial V} = \frac{4E_0}{V_e} \left[\left(\frac{V_e}{V} \right)^5 - \left(\frac{V_e}{V} \right)^3 \right]$$

One should note that $P(V_e) = 0$, which means that, with this model, the pressure to apply to the system at equilibrium is null.

$$\kappa(V) = - \frac{1}{V} \frac{\partial V}{\partial P} \Leftrightarrow \frac{1}{\kappa(V)} = -V \frac{\partial P}{\partial V} = \frac{4E_0}{V_e} \frac{V}{V_e} \left[5 \left(\frac{V_e}{V} \right)^6 - 3 \left(\frac{V_e}{V} \right)^4 \right] = \frac{4E_0}{V_e} \left[5 \left(\frac{V_e}{V} \right)^5 - 3 \left(\frac{V_e}{V} \right)^3 \right]$$

hence :

$$\mu(V) = \frac{\partial \frac{1}{\kappa}}{\partial P} = \frac{\partial \frac{1}{\kappa}}{\frac{\partial P}{\partial V}} = \frac{-\frac{4E_0}{V_e^2} \left[5 \cdot 5 \left(\frac{V_e}{V} \right)^6 - 3 \cdot 3 \left(\frac{V_e}{V} \right)^4 \right]}{-\frac{4E_0}{V_e^2} \left[5 \left(\frac{V_e}{V} \right)^6 - 3 \left(\frac{V_e}{V} \right)^4 \right]} = \frac{\left[25 \left(\frac{V_e}{V} \right)^2 - 9 \right]}{\left[5 \left(\frac{V_e}{V} \right)^2 - 3 \right]} \text{ and}$$

$$\boxed{\mu(V = V_e) = \mu(P = 0) = \frac{25 - 9}{5 - 3} = 8}$$

In the domain $0 < V \leq V_e$, $\mu(V)$ is an increasing function of V with $5 < \mu(V) \leq 8$.